# On ground state of atoms and molecules and some convexity issue

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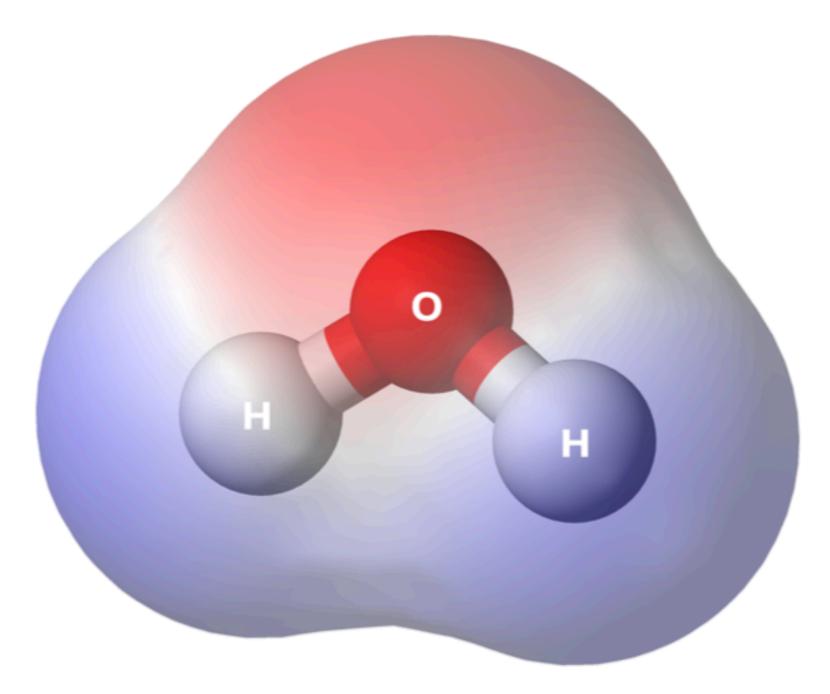
### **Plan for today lecture**

- 1. Density Functional Theory: from the quantum to the classical
- 2. Counter-example to the convexity conjecture

### Lot of calculus of variations

### **Some references:**

- E.H. Lieb, *Density functionals for Coulomb Systems*, International Journal of Quantum Chemistry, XXIV, 243-277, 1983.
- 2023.
- M. Lewin Coulomb and Riesz gases: The known and the unknown. J. Math. Phys., 63, p. 061101, 2022.

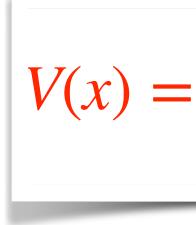


• M. Lewin, E.H. Lieb, R. Seiringer. Universal functionals in Density Functional Theory. Chapter 3 in Density Functional Theory — Modeling, Mathematical Analysis, Computational Methods, and Applications, edited by Eric Cances and Gero Friesecke, Springer,



# I. Density Functional Theory: from the quantum to the classical

### Schrödinger's equation for electrons in a molecule



$$H^{N}(\mathbf{V})\psi = E\psi, \quad H^{N}(\mathbf{V}) := -\frac{1}{2}\Delta_{\mathbb{R}^{3N}} + \sum_{j=1}^{N} \mathbf{V}(x_{j}) + \sum_{j< k} \frac{1}{|x_{j} - x_{k}|}$$

• M point nuclei of charges  $z_1, \ldots, z_M \in \mathbb{N}$  placed at  $R_1, \ldots, R_M \in \mathbb{R}^3$ 

$$= -\sum_{m=1}^{M} \frac{z_m}{|x - R_m|}$$

• N electrons: antisymmetric wave function  $\psi(x_1, \ldots, x_N)$  on  $(\mathbb{R}^3)^N$  with  $|\psi|^2 = 1$ 

**Problem:** essentially impossible to solve numerically!





### A variational principle and a double minimisation (Levy '82, Lieb '83)

Density of pure states

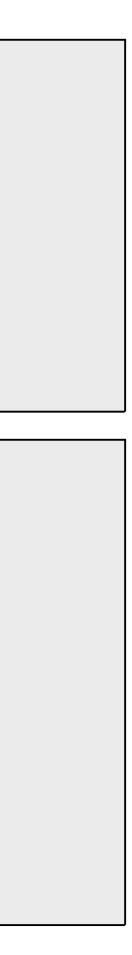
 $\langle \psi, H^N(V)\psi \rangle = \langle \psi, H^N(0)\rangle$ 

$$\rho_{\Psi}(x) = N \int_{\mathbb{R}^{3(N-1)}} |\Psi(x, x_2, \dots, x_N)|^2 dx_2 \cdots dx_N$$

Pure states

$$E[N, V] = \inf_{\psi} \langle \psi, H^{N}(V)\psi \rangle = \inf_{\rho} \left\{ \inf_{\psi \mid \rho_{\psi} = \rho} \langle \psi, H^{N}(0)\psi \rangle + \int_{\mathbb{R}^{3}} \rho(x)V(x)dx \right\}$$
$$F_{LL}[\rho] \text{ Levy-Lieb functional}$$

$$(D)\psi\rangle + \int_{\mathbb{R}^3} \rho_{\psi}(x) V(x) dx$$



### Some remarks

- New unknown  $\rho$  depends on only one variable  $x \in \mathbb{R}^3 \Rightarrow \underline{\text{no N-Particle space anymorel}}$
- Easy to show existence of  $F_{LL}$
- practice and no-explicit formulation in terms of  $\rho$
- It is not convex for
- **Density Functional Theory (DFT):**
- Understand better the true  $F_{LL}$ ;
- Replace it with  $F_{app}$  and then  $E[N, V] \approx \inf_{\rho} \left\{ F_{app}[\rho] + \int \rho V \right\}.$

•  $F_{LL}$  = "universal Levy-Lieb functional", very nonlinear and non local, impossible to compute in

### Let's try again: Legendre duality (Lieb '83)

#### Legendre transform and duality

 $V \mapsto E[N, V]$  is concave hence we can write

$$E[N, V] = \inf_{\rho} \left\{ F[\rho] + \int_{\mathbb{R}^3} \rho(x) V(x) dx \right\},$$
  
w-lsc functional defined by  
$$F[\rho] = \sup_{V} \left\{ E[N, V] - \int_{\mathbb{R}^3} \rho(x) V(x) dx \right\},$$

where  $F[\rho]$  is convex

$$F[\rho] + \int_{\mathbb{R}^3} \rho(x) V(x) dx \bigg\},$$
  
lefined by  
$$F[\rho] = \sup_{V} \bigg\{ E[N, V] - \int_{\mathbb{R}^3} \rho(x) V(x) dx \bigg\},$$

with  $\rho: \mathbb{R}^3 \to \mathbb{R}$  "variable dual to V"

#### Good news: *F* is convex!!

Bad news: Who is *F*?!?!



## Who is $\rho$ ? Who is $F[\rho]$ ?

**Density of mixed states** 

$$\operatorname{tr}(H^{N}(V)\Gamma) = \operatorname{tr}(H^{N}(0)\Gamma) + \int_{\mathbb{R}^{3}} \rho(x)V(x)dx$$

 $\rho_{\Gamma} = \sum n_j \rho_{\Psi_j},$ 

#### Rmk: $(V, \Gamma) \mapsto tr(H^N(V)\Gamma)$ is <u>linear</u> both in V and $\Gamma$

#### Mixed states: a double minimisation approach

$$E[N, V] = \inf_{\Gamma} \operatorname{tr}(H^{N}(V)\Gamma) = \inf_{\rho} \left\{ \inf_{\substack{\Gamma \mid \rho_{\Gamma} = \rho}} \operatorname{tr}(H^{N}(0)\Gamma) + \int_{\mathbb{R}^{3}} \rho(x)V(x)dx \right\}$$
  
$$F[\rho] \text{ Lieb functional}$$



$$\Gamma = \sum_{j} n_{j} |\Psi_{j}\rangle \langle \Psi_{j}|$$



### The universal functional $F[\rho]$

#### Theorem (Lieb '83)

The universal functional  $F[\rho]$ , satisfying the previous Legendre duality relations, is

It is finite if and only if  $\sqrt{\rho} \in H^1(\mathbb{R}^3)$ .

**Inf-sup argument** 

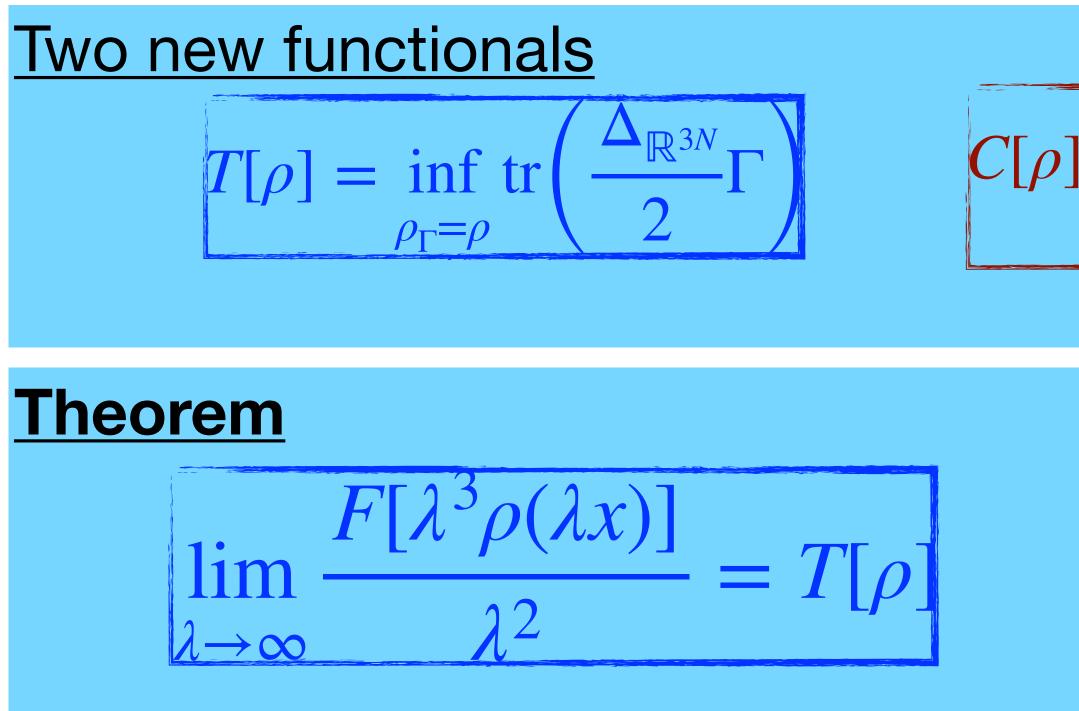
$$F[\rho] = \sup_{V} \left\{ E[N, V] - \int \rho V \right\} = \sup_{V} \inf_{\Gamma} \left\{ \operatorname{tr}(H^{N}(V)\Gamma) - \int \rho V \right\}$$
$$= \inf_{\Gamma} \sup_{V} \left\{ \operatorname{tr}(H^{N}(0)\Gamma) + \int (\rho_{\Gamma} - \rho)V) \right\} = \inf_{\Gamma} \left\{ \operatorname{tr}(H^{N}(0)\Gamma) + \sup_{V} \int (\rho_{\Gamma} - \rho)V \right\}$$
$$= + \infty \text{ unless } \rho_{\Gamma} = \rho$$

$$= \sup_{V} \left\{ E[N, V] - \int \rho V \right\} = \sup_{V} \inf_{\Gamma} \left\{ \operatorname{tr}(H^{N}(V)\Gamma) - \int \rho V \right\}$$
$$= \inf_{\Gamma} \sup_{V} \left\{ \operatorname{tr}(H^{N}(0)\Gamma) + \int (\rho_{\Gamma} - \rho)V) \right\} = \inf_{\Gamma} \left\{ \operatorname{tr}(H^{N}(0)\Gamma) + \sup_{V} \int (\rho_{\Gamma} - \rho)V \right\}$$
$$= + \infty \text{ unless } \rho_{\Gamma} = \rho$$

$$F[\rho] := \inf_{\rho_{\Gamma} = \rho} \operatorname{tr}(H^{N}(0)\Gamma)$$



### Low and high density regimes (aka some 1 - cv results)



#### **Rmk**:

- Convergence to  $T[\rho]$  rather easy (Lewin-Lieb-Seiringer '22)
- Bindini-De Pascale '18, Lewin '18)

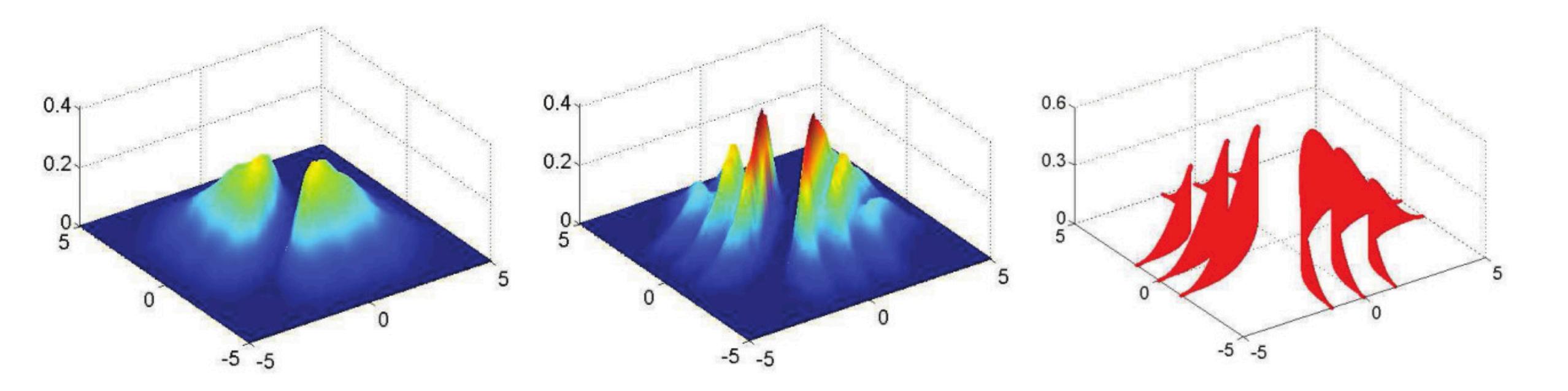
$$\int = \inf_{\mathbb{P} \text{ sym, } \rho_{\mathbb{P}} = \rho} \int_{\mathbb{R}^{3N}} \sum_{1 \le j \le k \le N} \frac{d\mathbb{P}(x_1, \dots, x_N)}{|x_j - x_k|}$$
$$\lim_{\lambda \to 0} \frac{F[\lambda^3 \rho(\lambda x)]}{\lambda} = C[\rho]$$

• Convergence to  $C[\rho]$  much more complicated due to the lack of regularity of classical problem (Cotar-Friesecke-Klüppelberg '13-'18,





### The low-density limit is (very) singular



Pair density for 
$$\lambda = 1$$
,  $\lambda = 0$   
 $\rho(x) = \frac{2}{5}(1 + \frac{1}{5})$ 

 $\lambda = 0.1$  and  $\lambda \approx 0$  in 1D, with N = 4 and  $\frac{2}{5}(1 + \cos(\pi x/5))\chi_{x \in [-5,5]}(x)$ (Chen-Friesecke '15)

### Some remarks

- Existence for  $F_{LL}[\rho]$  and  $F[\rho]$ : it follows by using the direct method fo calculus of variations. Notice that the mass constraint (e.g.  $||\psi|^2 = 1$ ) helps to get some compactness.
- $\rho \mapsto F_{LL}[\rho]$  is non convex so the Legendre duality fails.
- Relation between  $F_{LL}$  and F:

$$F[\rho] = \inf\left\{\sum_{j} n_j F_{LL}[\rho_j] \mid \sum_{j} n_j \rho_j = \rho, \sum_{j} n_j = 1, n_j \ge 0\right\}$$

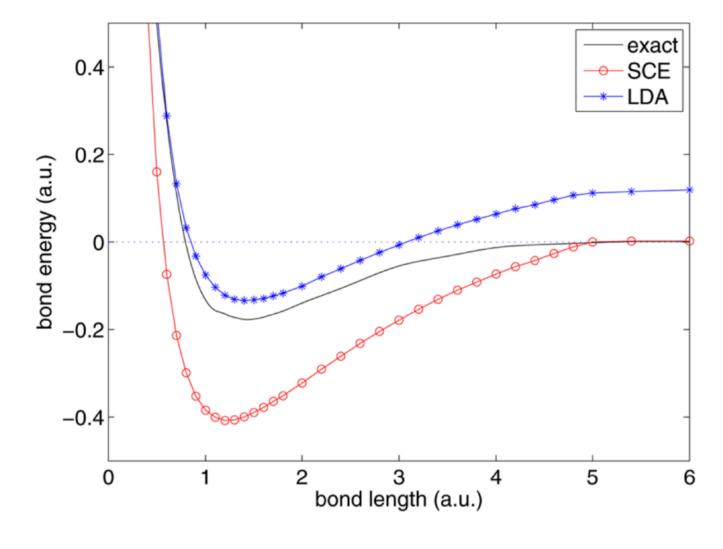
F is actually the convex hull of  $F_{LL}$ .

I totally ignored the 3 hours speech on the functional spaces but

$$\rho \in \left\{ \int \rho = N, \sqrt{\rho} \in H^1(\mathbb{R}^3) \right\} \text{ and } V \in L^{3/2}(\mathbb{R}^3) + L^{\infty}(\mathbb{R}^3)$$

- $C[\rho]$  is a (multi-marginal) optimal transport problem.
- In computational chemistry  $F[\rho] \approx T[\rho] + C[\rho]$

 $\mathbb{R}^3$ ) (Lieb '83).

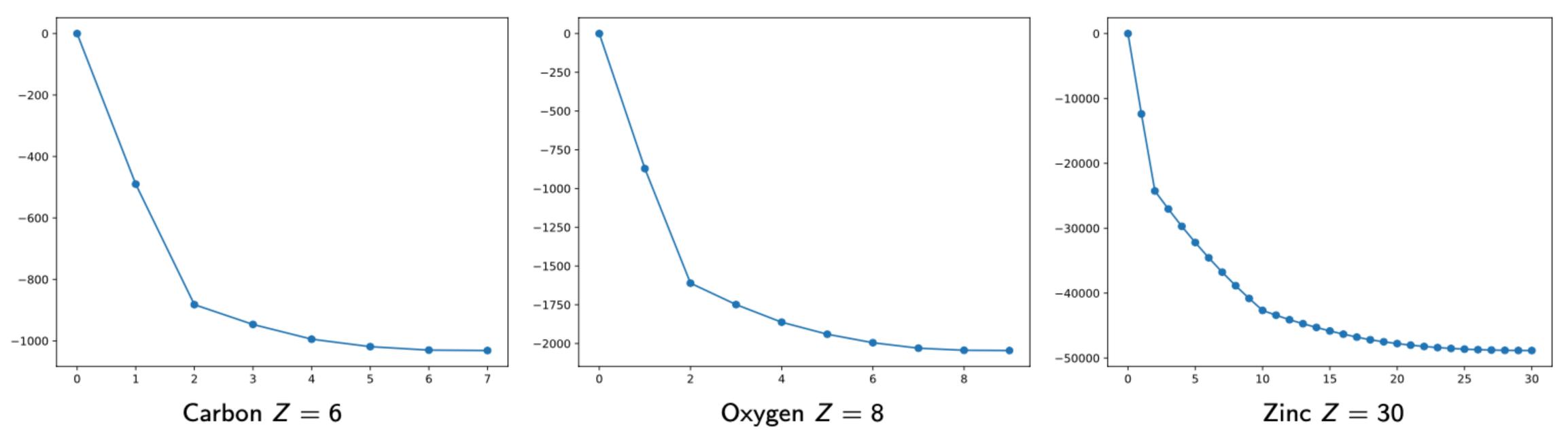


 $H_2$  dissociation Chen-Friesecke-Mendl '14



## II. A convexity conjecture in Quantum Chemistry

### Experimental energies of an atom with Z protons, in terms of the number **Nof electrons:**



**Rmk #1**: monotone (additional electron stays if energy decreases, or escapes to infinity) **Rmk #2: convex** (valence electrons less tightly bound than core electrons) **Rmk #3**: constant for  $N \ge Z + 1$  or Z + 2 (nucleus cannot bind too many electrons)

Electrons described by Schrödinger's equation!

- #1 easy to prove
- #2, #3 very hard to prove  $\Rightarrow$  "ionisation conjecture"

Data from Wikipedia, NIST



# **Ionization conjecture #2**

#### <u>Conjecture (convexity-in-N)</u>

#### For $V \in ?$ , the map $N \mapsto E[N, V]$ is convex, which means (with E[0, V] = 0)

• Perdew-Parr-Levy-Balduz (1982) and Parr-Yang (1994)

suggest conjecture true for all  $V(x) = -\sum_{m=1}^{M} \frac{z_m}{|x - R_m|}, z_m \in \mathbb{N}$ 

 $^{\circ}$  Lieb (1983) stated the conjecture for all V

While it has been conjectured that E(N, v) is convex in N (for all v) in the case of Coulomb repulsion, this has never been proved. It has not even been proved that  $E(3, v) + E(1, v) \ge 2E(2, v)$ .

• Simon (1984) Included the conjecture for atoms in a famous list of 15 open problems

#### $E[N, V] - E[N - 1, V] \le E[N + 1, V] - E[N, V], \quad \forall N \in \mathbb{N}$

$$E^{0}(N+1) - E^{0}(N) \ge E^{0}(N) - E^{0}(N-1)$$
(4.1.14)

or

$$I(N+1) \ge I(N)$$
 (4.1.15)

where I(N) is the ionization potential of the N-electron ground state. Equation (4.1.15) states that successive ionization potentials are not decreasing (for fixed external potential).

For atoms and molecules, no counterexample is known to (4.1.15), although a first-principles proof has never been given. As examples, in

**Problem 10A** (Monotonicity of the Ionization Energy). Prove that

 $(\Delta E) (N-1, Z) \ge (\Delta E) (N, Z)$ 

for all N, Z.

This is just the fact, almost obvious, that it takes more energy to remove inner electrons than outer ones. Since in removing electron (N-1) there is one fewer electron to repel, and since the Pauli principle only makes things better this should be true. It seems to be remarkably difficult to prove.



### A counter-example with nuclei of fractional charges

#### **Theorem**

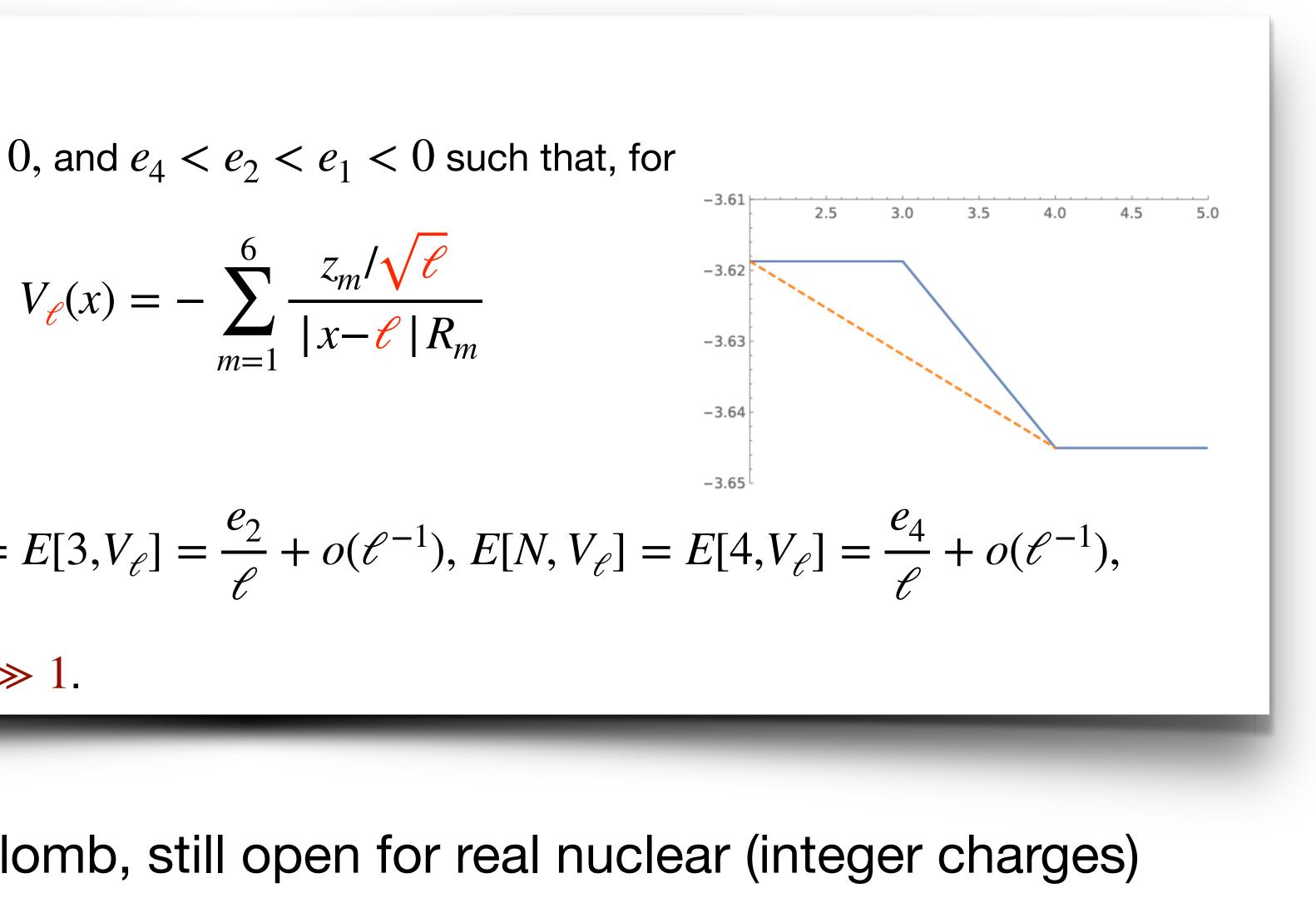
There exists  $R_1, ..., R_6 \in \mathbb{R}^3, z_1, ..., z_6 > 0$ , and  $e_4 < e_2 < e_1 < 0$  such that, for

We have for all  $N \ge 5$ 

$$E[1, V_{\ell}] = \frac{e_1}{\ell} + o(\ell^{-1}), E[2, V_{\ell}] = E[3, V_{\ell}] = E[3,$$

and hence **convexity fails at** N = 3 for  $\ell \gg 1$ .

- Follows from calculus of variations arguments for classical electrons



First counter-example for Coulomb, still open for real nuclear (integer charges)

# **Back to classical** N-body problem **Classical minimal energy** $e[N, V] = \inf_{x_1, \dots, x_N} \left( \sum_{i=1}^N \sum_{i=1}^N \right)^{-1}$

#### <u>Conjecture (convexity-in-N)</u>

For  $V \in ?$ , the map  $N \mapsto e[N, V]$  is convex

Quantum conjecture for all nice-enough  $V \Rightarrow$  classical one since

 $E[N, \varepsilon V(\varepsilon \cdot)]$ lim  $\varepsilon \rightarrow 0$  $\mathcal{E}$ 

$$\frac{V(x_j)}{1 \le j \le k \le N} \frac{1}{|x_j - x_k|}$$

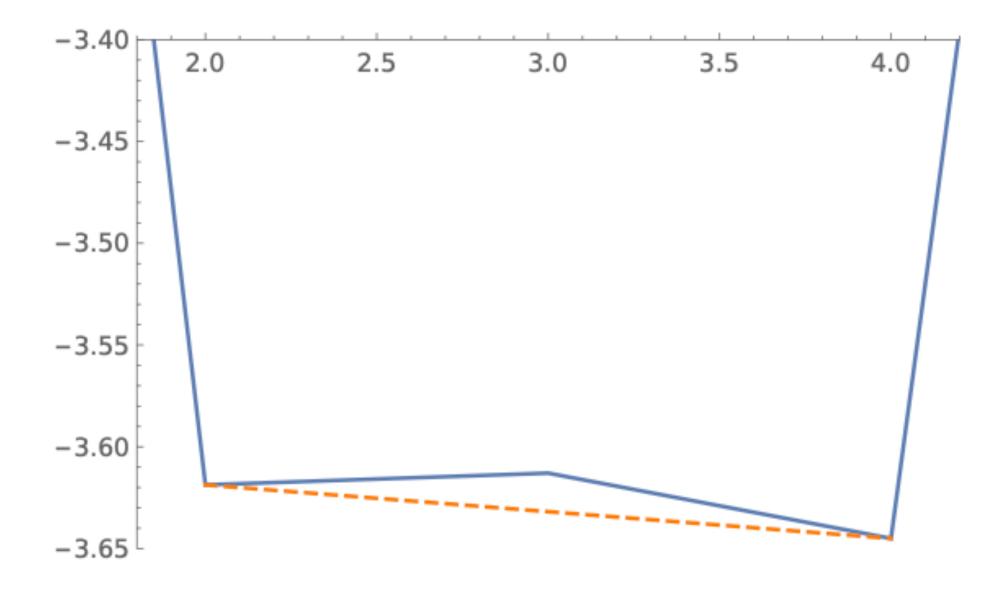
# = e[N, V]



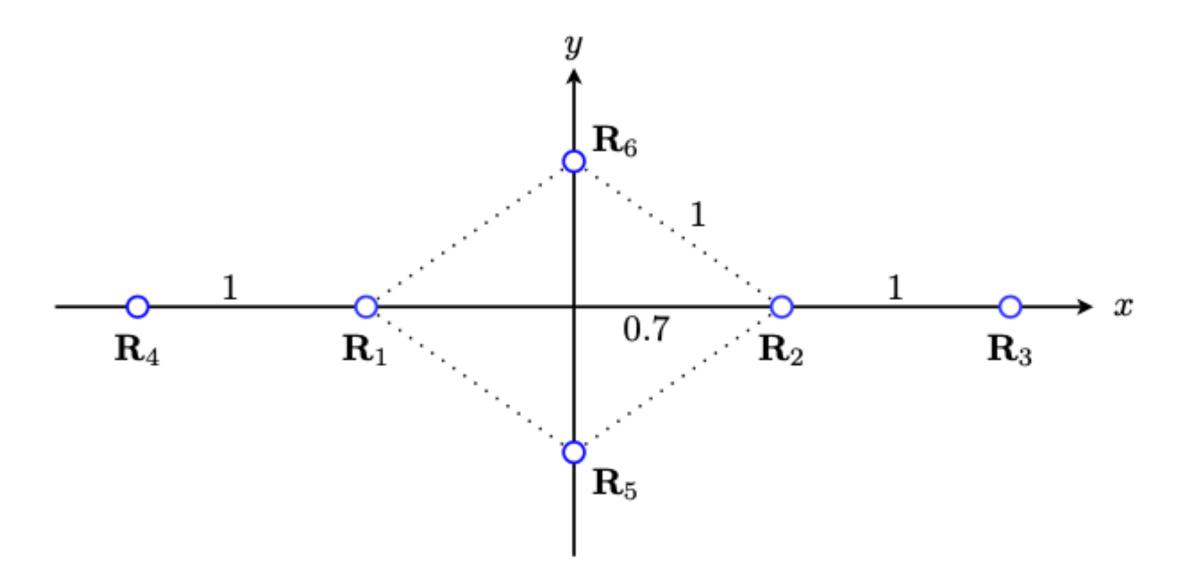
#### Classical counter-example

$$V(x) = \begin{cases} v_m & \text{if } x = R_m \\ +\infty & \text{if } x \notin \{R_1, \dots, R_6\} \end{cases}$$

$$\begin{cases} v_1 = v_2 = -2.1665 \\ v_3 = v_4 = -1.4109 \\ v_5 = v_6 = -1.9934 \end{cases}$$



Ν	$E_{ m cl}^{N}[V] pprox$	minimizer
1	-2.1665	$\mathbf{R}_1$
2	-3.6187	$\mathbf{R}_1, \mathbf{R}_2$
3	-3.6129	$\mathbf{R}_4, \mathbf{R}_5, \mathbf{R}_6$
4	-3.6450	$R_3,, R_6$
5	-2.3949	$R_2,, R_6$
6	-0.4304	$R_1,, R_6$



**Variational formulation for the convex hull of** 
$$N \mapsto e(N, V)$$
  
 $e(N, V) = \inf E_N^V = \inf_{\mathbb{P} \text{ symm}} \int_{\mathbb{R}^{3N}} E_N^V d\mathbb{P}, \quad E_N^V(x_1, \dots, x_N) = \sum_{j=1}^N V(x_j) + \sum_{1 \le j \le k \le N} \frac{1}{|x_j - x_k|}$ 

**Grand-canonical: make** *N* **random** 

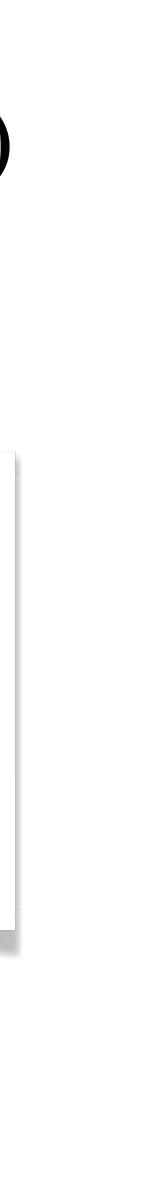
$$e_{GC}(\lambda, \mathbf{V}) = \inf \left\{ \sum_{n \ge 1} p_n e(n, \mathbf{V}) \mid \sum_n p_n = 1, \sum_n n p_n = \lambda \right\}$$

That is,  $\lambda \in \mathbb{R}_+ \mapsto e_{GC}(\lambda, V)$  is the **convex hull** of  $N \mapsto e(N, V)$ 

**Example:** for previous counter-example V, better to have 2 and 4 particles each with proba 1/2 instead of 3.

Conjecture true for one V and all  $N \ge 1$ 

$$\iff e_{GC}(N, V) = e(N, V), \forall N \ge 1$$



### Legendre transforms (and duality) Let's look at the Legendre transforms of e[N, V] and $e_{GC}[N, V]!$

Multi-marginal optimal transport with Coulomb cost

$$C[\rho] = \inf_{\mathbb{P} \text{ sym, } \rho_{\mathbb{P}} = \rho} \int_{\mathbb{R}^{3N}} \sum_{1 \le j \le k \le N} \frac{d\mathbb{P}(x_1, \dots, x_N)}{|x_j - x_k|}$$

then

$$C[\rho] = \sup_{V} \left\{ e[N, V] - \int \rho V \right\},$$

**Grand-canonical optimal transport with Coulomb cost** 

$$C_{GC}[\rho] = \inf\left\{\sum_{n\geq 1} p_n C[\rho_n] \mid \sum_n p_n = 1, \sum_n p_n \rho_n = \rho\right\}$$

Conjecture true for all V and all  $N \ge 1 \iff C_{GC}[\rho] = C[\rho], \forall \rho$  with  $\rho = N$ 

$$e[N,V] = \inf_{\int \rho = N} \left\{ C[\rho] + \int \rho V \right\}$$

### Support for the grand canonical aka $C[\rho] \stackrel{\prime}{=} C_{GC}[\rho]$ For $p = (p_n)_{n>0}$ we call supp $(p) = \{n \mid p_n \neq 0\}$ its support in n

<u>Theorem (support in *n*)</u> Let  $\rho \ge 0$  with  $N = \rho(\mathbb{R}^3) \in \mathbb{N}$  and  $C_{GC}[\rho] \le C[\rho] < +\infty$ . Any optimiser for  $C_{GC}[\rho]$  satisfies  $\sup(p) \begin{cases} = \{N\} & \text{if } N \in \{0,1,2\}, \text{ hence } C_{GC}[\rho] = C[\rho] \\ \subset [N - \frac{1}{2}\sqrt{8N + 9} + \frac{3}{2}, N + \frac{1}{2}\sqrt{8N - 7} - \frac{1}{2}] & \text{if } N \ge 3 \end{cases}$ 

<u>Theorem (counter-example)</u>

There exists a  $\rho$  with  $\rho(\mathbb{R}^3) = 3$  such that  $\operatorname{supp}(p) = \{2,4\}$ , hence  $C_{GC}[\rho] < C[\rho]$ . Moreover, for every  $k \ge 1$ , There exists  $\rho^{(k)}$  with  $\rho^{(k)}(\mathbb{R}^3) = \frac{6^k}{2}$  such that  $\operatorname{supp}(p^{(k)}) = \left\{ \frac{6^k - 2^k}{2}, \frac{6^k + 2^k}{2} \right\}$ 

### $\Rightarrow$ convexity in N conjecture cannot hold!

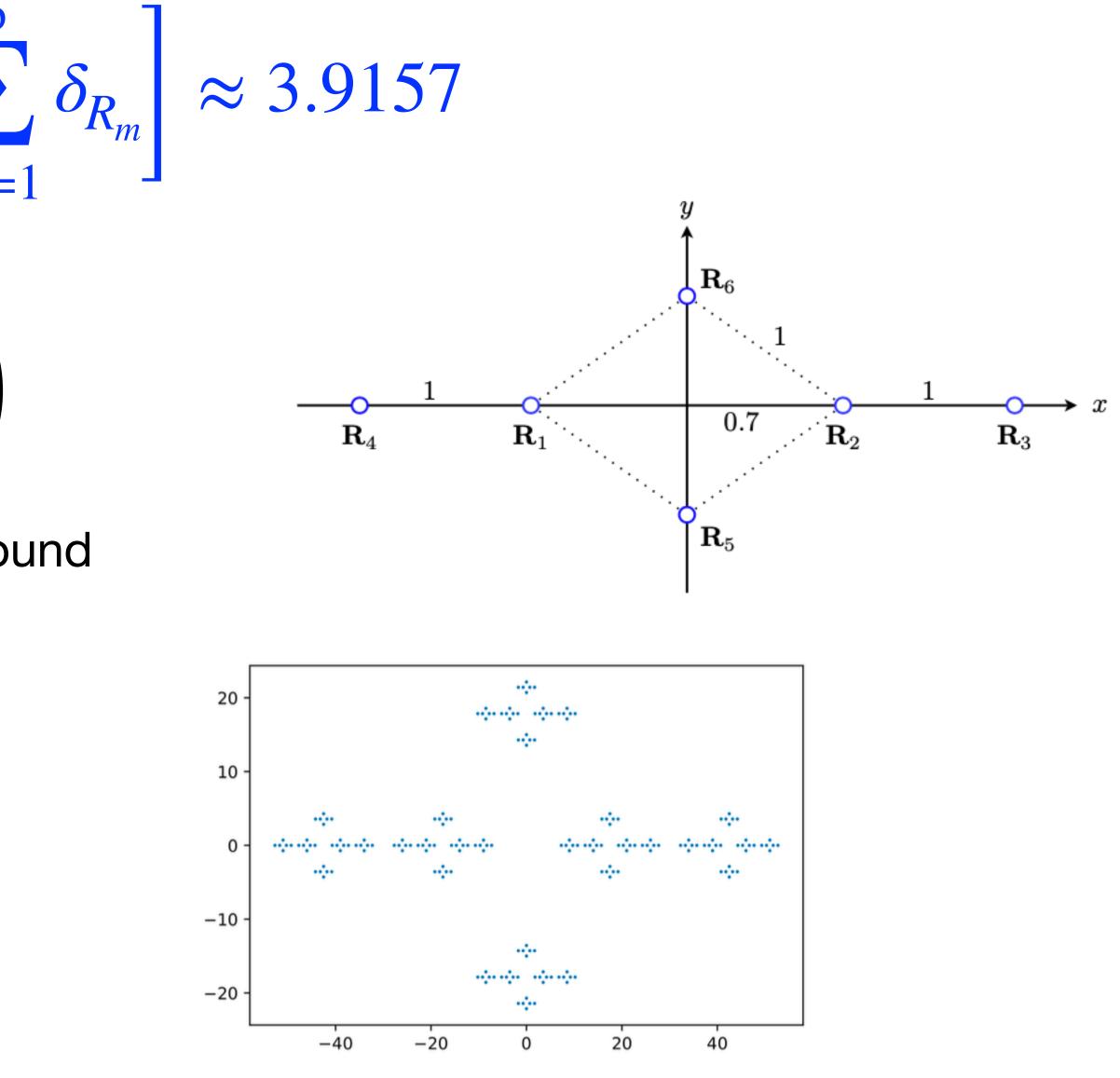


• For the 6 points  $R_1, \ldots, R_6$  below-right we have

$$3.8778 \approx C_{GC} \left[ \frac{1}{2} \sum_{m=1}^{6} \delta_{R_m} \right] < C \left[ \frac{1}{2} \sum_{m=1}^{6} \delta_{R_m} \right]$$

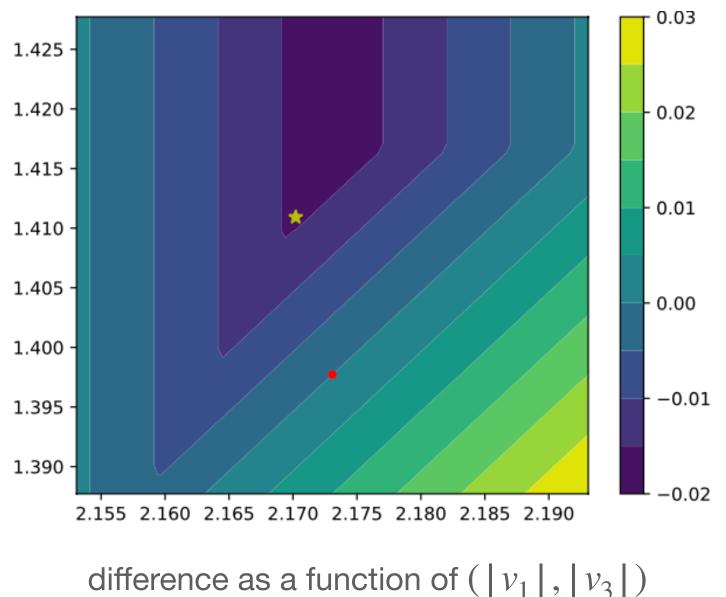
- With the optimiser  $\mathbb{P} = \frac{1}{2} \left( \delta_{R_1} \otimes \delta_{R_2} + \delta_{R_3} \otimes \ldots \otimes \delta_{R_6} \right)$
- Repeating this pattern at different scales we found

$$\operatorname{supp}(p^{(k)}) = \left\{ \frac{6^k - 2^k}{2}, \frac{6^k + 2^k}{2} \right\}$$



### Finding the potential V• To find the potential V, we first solve the dual problem to $C_{GC}$ bu we get

- - $e[2,V_{GC}] = e[3,V_{GC}] = e[4,V_{GC}]$
- Idea: minimize  $V \mapsto (e[2,V] + e[4,V])/2 e[3,V]$  in a neighbourhood of  $V_{GC}$  to get V.



# Conclusion

- convexity-in-N conjecture wrong for general Coulomb potentials
- still open for atoms and molecules
- experiments and numerics say it is true for atoms, but no (mathematical) intuition why
- very helpful to work with tools from calculus of variations (Legendre transform)
- Schrödinger equations is 100 years old in 2025 but still poses many interesting mathematical questions with large impact in applications

### Not yet the end (Exam on 17/01/2025 room 0A7) • Existence: use the direct method of calculus of variations so look for (1) compactness and (2) lower semi-continuity.

- Uniqueness follows from strict convexity of the functional.
- Euler-Lagrange equations hold usually in the sense of distributions: prove that they hols in a stringer sense demands some additional work (remember the one dimensional case).
- Legendre transform and duality can help to understand better the problem, that is the properties of the minimiser (regularity).
- •Variational convergence, aka  $\Gamma$  convergence, helps to understand better a problem if we approximate it with a suitable optimization problem.

Happy 2025! And good luck (especially if you go for a Ph.D. whatever the topic )!

